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# Determination of long range dipolar couplings in ferroelectric liquid crystals 

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#### Abstract

Long range dipolar coupling constants have been determined in three ferroelectric liquid crystals in their racemic forms using ${ }^{13} \mathrm{C}$ NMR. Two of these liquid crystals are esters of $\alpha$-chloroacids and 4 -octyloxy-4'-hydroxybiphenyl, and have a very large spontaneous polarization in the smectic C* phase. The strategy used in the present study is the observation and measurement of ${ }^{2} \mathrm{H}^{13} \mathrm{C}$ splittings in the ${ }^{13} \mathrm{C}$ spectra of monodeuterated compounds The order parameters were calculated from the 1D spectra, and some of the coupling constants are compared with the ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ coupling constants previously obtained from 2D experiments. In addition, the deuterium quadrupole splitting of these compounds was determined from their ${ }^{2} \mathrm{H}$ NMR spectra. The experiments were carried out over the whole mesomorphic ranges of the liquid crystals, covering the smectic A and smectic C phases.


## 1. Introduction

Ferroelectric liquid crystals (FLCs) have attracted considerable attention since their discovery in 1975 [1]. These materials form a chiral smectic C (SmC*) phase which exhibits spontaneous electrical polarization. Hence, considerable effort has been spent on the search for FLC molecules with large spontaneous polarization. An important class of FLCs is formed by the esters of optically active $\alpha$-halocarboxylic acids with 4 -alkoxy-$4^{\prime}$-hydroxybiphenyls [2-4]. These compounds exhibit a chiral smectic C phase with very large spontaneous polarization $\mathbf{P}_{\mathrm{S}}$ (up to $300 \mathrm{nCcm}^{-2}$ ).

A variety of techniques has been used to study molecular motion in liquid crystalline phases. Among them, NMR spectroscopy has been used very successfully to give detailed information on molecular geometry, orientational ordering and dynamics of molecular rotations [5]. Orientational ordering is one of the most fundamental characteristics of LC molecules. It has been shown that the 2D NMR method of separated local field (SLF) spectroscopy [6] and proton encoded local field (PELF) spectroscopy [7, 8], especially in combination with off-magic-angle spinning (OMAS), is very effective in studying the ordering of liquid crystals. The

[^0]order parameters of several FLCs have been determined by the use of ${ }^{13} \mathrm{C}$ and ${ }^{19} \mathrm{~F}$ NMR [9-13].

Molecular dynamics is perhaps the most powerful theoretical method for conformational studies of LCs. The criterion of its success is a comparison between the calculated order parameters of molecular segments and the corresponding experimental data. A more stringent test for the molecular dynamics calculations is a comparison between calculated and experimental long range dipolar coupling constants, which can be determined by using several different NMR techniques [14-19]. We have found that the use of monodeuterated LCs [14, 19] to observe ${ }^{2} \mathrm{H}-{ }^{13} \mathrm{C}$ splittings in the ${ }^{13} \mathrm{C}$ spectra gives very rich information [19]; the application of this approach to determine some long range dipolar coupling constants in three racemic ferroelectric liquid crystals is reported in this work. The compounds studied are $4^{\prime}$-( 4 -methyl-2-chloro-2-deuteropentanoyloxy)-4-octyloxybiphenyl (4M2C2DPOOB), 4'-(3-methyl-2-chloro-2-deutero-pentanoyloxy)-4 -octyloxybiphen yl (3M2C2DPOOB) and 4'-(4-hexylbenzoyloxy)-4- (1-methyl-1-deuteroheptyloxy) biphenyl (HBO1M1DHOB). To simplify the spectral analysis, we also studied the protonated liquid crystals, 4'-(4-methyl-2-chloropentanoyloxy)-4-octyloxybipheny 1 (4M2CPOOB), 4'-(3-methyl-2-chloropentanoyloxy)-4-octyloxybi phenyl (3M2CPOOB) and 4'-(4-hexylbenz oyl-oxy)-4-(1-methylheptyloxy)biphenyl (HBO1MHOB) by ${ }^{13} \mathrm{C}$ NMR. The ${ }^{2} \mathrm{H}$ NMR spectra of these compounds also provide the quadrupolar splittings of the deuterons. In addition, the one- and two-bond ${ }^{1} \mathrm{H}^{13} \mathrm{C}$ dipolar


Scheme 1.

coupling constants of 3 M 2 CPOOB are compared with those previously determined in this laboratory using the 2D SLF method [9].

## 2. Experimental

3M2CPOOB was synthesized according to scheme 1 as described in the literature [2, 3]. Its isomer 4M2CPOOB was prepared similarly. The transition temperatures of the purified 3 M 2 CPOOB are: $63.5^{\circ} \mathrm{C}$ from isotropic liquid to SmA phase, $46.5^{\circ} \mathrm{C}$ from SmA to SmC phase, $30.0^{\circ} \mathrm{C}$ from SmC to SmG phase and $20.0^{\circ} \mathrm{C}$ from SmG to solid. The transition temperatures of the purified 4 M 2 CPOOB are: $68.6^{\circ} \mathrm{C}$ from isotropic liquid to SmA , $56.5^{\circ} \mathrm{C}$ from SmA to SmC and $52.5^{\circ} \mathrm{C}$ from SmC to solid.

The monodeuterated liquid crystal 3M2C2DPOOB was synthesized according to scheme 2. 2-Deutero-2-hydroxy-3-methylpentanoic acid was prepared by reduction of the corresponding ketone using $\mathrm{NaBD}_{4}$ in methanol. Then, the acid chloride was obtained by
action of thionyl chloride on the acid. Finally, the acid chloride was esterified with 4-octyloxy-4'-hydroxybiphenyl and pyridine in dry chloroform. The final compound was recrystallized from hexane. Its isomer 4M2C2DPOOB was prepared similarly.

The monodeuterated liquid crystal HBO1M1DHOB was synthesized according to scheme 3. 2-Octanone was reduced with $\mathrm{NaBD}_{4}$ in methanol to give the corresponding alcohol. This alcohol was then transformed to the chloride compound with thionyl chloride. The next step was the mono-alkylation of the chloride with $4,4^{\prime}$-biphenol using $\mathrm{NaHCO}_{3}$ as base in a mixture of dioxan/PEG (polyethylene glycol) as solvent. Finally, the phenol was esterified with $4^{\prime}$-hexylbenzoyl chloride in pyridine and dry chloroform. The final compound was recrystallized from ethanol. The protonated analogue was prepared similarly. The transition temperatures of the purified HBO1M1DHOB are: $85^{\circ} \mathrm{C}$ from isotropic liquid to nematic ( N ) phase, $51.8^{\circ} \mathrm{C}$ from N to SmC phase and $40^{\circ} \mathrm{C}$ from SmC to solid.



Scheme 2.





Scheme 3.


The $1 \mathrm{D}{ }^{13} \mathrm{C}$ and ${ }^{2} \mathrm{H}$ NMR experiments were performed using a Varian UNITY/INOVA 400 NMR spectrometer. The sample was placed in a standard 5 mm tube and rotated slowly along the magnetic field so that the director aligned parallel to the field. The temperature calibration was made by observing the nematic to isotropic transition for the compounds studied.

## 3. Results and discussion

3.1. ${ }^{13} \mathrm{C}$ NMR study of monodeuterated liquid crystals

In an early ${ }^{13} \mathrm{C}$ NMR study of a monodeuterated liquid crystal, long range ${ }^{2} \mathrm{H}^{13} \mathrm{C}$ splittings were not well resolved because of inefficient ${ }^{1} \mathrm{H}$ decoupling [14]. For
our recent work [19] and the present study, we used a very efficient broadband decoupling sequence for liquid crystals developed in our laborator y [20], which removes all ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ couplings to yield sharp ${ }^{13} \mathrm{C}$ peaks and requires only moderate decoupler power ( $\gamma B_{2} / 2 \pi<20 \mathrm{kHz}$ ). The ${ }^{13} \mathrm{C}$ spectra of the monodeuterated liquid crystals $4 \mathrm{M} 2 \mathrm{C} 2 \mathrm{DPOOB}, 3 \mathrm{M} 2 \mathrm{C} 2 \mathrm{DPOOB}$, and HOB1M1DHOB and their protonated analogues obtained by the use of this decoupling sequence (SPINAL-64) are shown in figures 1,2 , and 3 , respectively.

For 4M2C2DPOOB, the coupling of the ${ }^{2} \mathrm{H}$ nucleus with five types of carbon (the carbonyl, the $3^{\prime}$ - and $4^{\prime}$-carbons of the nearby phenyl ring, and the carbons $\alpha^{\prime}$ and $\beta^{\prime}$ of the aliphatic chain) is resolvable (figure 1 ).


Bottom



Figure 1. ${ }^{13} \mathrm{C}$ NMR spectra of 4 M 2 CPOOB (top) and its monodeuterated analogue 4 M 2 C 2 DPOOB (bottom) at 100.58 MHz and $65^{\circ} \mathrm{C}$. The spectra were obtained with cross polarization, using the SPINAL- 64 broadband ${ }^{1} \mathrm{H}$ decoupling sequence with $\gamma B_{2} / 2 \pi=18 \mathrm{kHz}$. The first set of vertical bars shows the peak positions of the non-deuterated homologue; the second set of vertical bars shows the positions of the outer peaks for each resolvable triplet due to ${ }^{1} \mathrm{H}$ splitting.


Figure 2. ${ }^{13} \mathrm{C}$ NMR spectra of 3 M 2 CPOOB (top) and its monodeuterated analogue 3 M 2 C 2 DPOOB (bottom) at 100.58 MHz and $37^{\circ} \mathrm{C}$. Two sets of diastereo isomers, a $[(2 S, 3 S)$ and $(2 R, 3 R)]$ pair and a $[(2 S, 3 R)$ and $(2 R, 3 S)]$ pair, were present in each compound. The spectra were obtained with cross polarization, using the SPINAL-64 broadband ${ }^{1} \mathrm{H}$ decoupling sequence with $\gamma B_{2} / 2 \pi=18 \mathrm{kHz}$. The first set of vertical bars shows the peak positions of the non-deuterated homologue; the second set of vertical bars shows the positions of the outer peaks for each resolvable triplet due to ${ }^{2} \mathrm{H}$ splitting.

On the aromatic side, the $3^{\prime}$ carbon is separated from the deuterium atom by five bonds. In the aliphatic region, many peaks overlap with each other because of the large number of aliphatic carbons in the molecule. Therefore, although the spectrum of the monodeuterated compound shows a new peak at 26 ppm , unambiguous ${ }^{2} \mathrm{H}-{ }^{1} \mathrm{C}$ dopolar couplings could be obtained only for the $\alpha^{\prime}$ and $\beta^{\prime}$ carbons, which are one and two bonds away from the deuterium, respectively.

For 3M2CPOOB, the ${ }^{13} \mathrm{C}$ peaks for the two sets of diastereo isomers are distinguishable (figure 2). Seven sets of ${ }^{2} \mathrm{H}_{-}{ }^{13} \mathrm{C}$ splittings are observed for 3 M 2 C 2 DPOOB due to couplings of the two sets of isomers with the $\alpha^{\prime}$, $\beta^{\prime}, \omega^{\prime}$, and $\omega^{\prime \prime}$ carbons. Further discussion on this point will be made later. In the aromatic region, the peaks are broad and the ${ }^{2} \mathrm{H}-{ }^{13} \mathrm{C}$ couplings could not be resolved.

For the liquid crystal HBO1M1DHOB, the coupling of the ${ }^{2} \mathrm{H}$ nucleus with five types of carbons (the 12 - and 11 - and 9 -carbons of the nearby phenyl ring, and the carbons $\alpha$ and $\beta$ of the aliphatic chain) was observed (figure 3). Splittings for the 10 -carbon, $\alpha-\mathrm{CH}_{3}$, and some other carbon atoms in the aliphatic region are clearly
visible, but we could not determinate the corresponding values of the dipolar coupling constants because of overlapping peaks. The 9 -carbon is separated from the deuterium atom by six bonds; although the splitting is not obvious in the spectrum, its value could be determined by spectral deconvolution.

The ${ }^{2} \mathrm{H}-{ }^{13} \mathrm{C}$ dipolar coupling constant $(D)$ is related to the observed splitting $\Delta v$ by

$$
\begin{equation*}
\Delta v=2 D+J \tag{1}
\end{equation*}
$$

where $J$ is the scalar coupling constant. The signs of $\Delta v$ cannot be determined, and are assumed to be positive based on geometric considerations. The $J$ values were determined from spectra in the isotropic phase, and only the directly bonded pairs were non-zero. For some of the peaks, spectral deconvolution was needed to obtain the value of $\Delta v$. A complete analysis of the results for the three monodeuterate $d$ compounds has been made over their entire liquid crystalline ranges. For comparison with previous data, the ${ }^{2} \mathrm{H}^{13} \mathrm{C}$ dipolar coupling constants are converted to ${ }^{1} \mathrm{H}^{-13} \mathrm{C}$ dipolar coupling constants




Figure 3. ${ }^{13} \mathrm{C}$ NMR spectra of HBO1MHOB (top) and its monodeuterated analogue HBO1M1DHOB (bottom) at 100.58 MHz and $48^{\circ} \mathrm{C}$. The spectra were obtained with cross polarization, using the SPINAL- 64 broadband ${ }^{1} \mathrm{H}$ decoupling sequence with $\gamma B_{2} / 2 \pi=18 \mathrm{kHz}$. The first set of vertical bars shows the peak positions of the non-deuterated homologue; the second set of vertical bars shows the positions of the outer peaks for each resolvable triplet due to ${ }^{2} \mathrm{H}$ splitting.
using the equation:

$$
\begin{equation*}
D_{\mathrm{CH}}=\left(\gamma_{\mathrm{H}} / \gamma_{\mathrm{D}}\right) D_{\mathrm{CD}}=6.5142 D_{\mathrm{CD}} \tag{2}
\end{equation*}
$$

where $\gamma_{\mathrm{H}}$ and $\gamma_{\mathrm{D}}$ are the magnetogyric ratios of proton and deuteron, respectively. The data for the compounds $4 \mathrm{M} 2 \mathrm{CPOOB}, 3 \mathrm{M} 2 \mathrm{CPOOB}$ and HBO1MHOB are presented in tables 1,2 , and 3 , respectively.

Most of the dipolar coupling constants listed in the tables decrease with the increase of temperature, reflecting the decreased ordering of the LC molecules. However, a few of them (e.g. $\alpha^{\prime}-\beta^{\prime}, \alpha^{\prime}-\omega^{\prime}$, and $\alpha^{\prime}-\omega^{\prime \prime}$ in table 2) do not follow this trend because these values depend on the average of many conformations. Although it is difficult to use the long range dipolar coupling constants to extract order parameters, they serve as a useful database for the study of conformational properties of LCs and as stringent criteria for judging the success of molecular dynamics simulations.

Table 1. ${ }^{1} \mathrm{H}_{-}{ }^{13} \mathrm{C}$ dipolar coupling constants for 4 M 2 CPOOB at different temperatures.

| $T /{ }^{\circ} \mathrm{C}$ | $\alpha^{\prime}-\mathrm{CO}$ | $\alpha^{\prime}-3$ | $\alpha^{\prime}-4$ | $\alpha^{\prime}-\alpha^{\prime}$ | $\alpha^{\prime}-\beta^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 53 | - | - | - | 7574 | 795 |
| 54 | - | - | - | 7446 | 772 |
| 55 | - | - | - | 6954 | 762 |
| 56 | 1362 | - | - | 6944 | 759 |
| 57 | 1316 | - | - | 6931 | 756 |
| 58 | 1287 | - | - | 6908 | 755 |
| 59 | 1257 | - | - | 6876 | 747 |
| 60 | 1251 | - | - | 6840 | 745 |
| 61 | 1228 | - | - | 6804 | 743 |
| 62 | 1189 | 180 | - | 6775 | 743 |
| 63 | 1124 | 175 | - | 6746 | 739 |
| 64 | 1124 | 171 | 168 | 6711 | 730 |
| 65 | 1120 | 170 | 160 | 6697 | 728 |
| 66 | 1117 | 169 | 152 | 6671 | 726 |
| 67 | 1114 | 168 | 140 | 6633 | 726 |
| 68 | 1107 | 165 | 132 | 6602 | 723 |
|  |  |  |  |  |  |

Table 2. ${ }^{1} \mathrm{H}_{-}{ }^{13} \mathrm{C}$ dipolar coupling constants for 3 M 2 CPOOB at different temperatures. The values listed correspond to two sets of diastereo isomers, a $[(2 S, 3 S)$ and $(2 R, 3 R)]$ pair and a $[(2 S, 3 R)$ and $(2 R, 3 S)]$ pair present in each compound. One set of $\alpha^{\prime}-\omega^{\prime}$ coupling could not be determined because of peak overlapping.

| $T /{ }^{\circ} \mathrm{C}$ | $\alpha^{\prime}-\alpha^{\prime}$ | $\alpha^{\prime}-\alpha^{\prime}$ | $\alpha^{\prime}-\beta^{\prime}$ | $\alpha^{\prime}-\beta^{\prime}$ | $\alpha^{\prime}-\omega^{\prime \prime}$ | $\alpha^{\prime}-\omega^{\prime \prime}$ | $\alpha^{\prime}-\omega^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 35 | 7635 | 7713 | 933 | 827 | 202 | 189 | 117 |
| 37 | 7517 | 7654 | 886 | 827 | 195 | 182 | 111 |
| 39 | 7433 | 7537 | 899 | 814 | 195 | 163 | 104 |
| 41 | 7289 | 7413 | 860 | 788 | 182 | 156 | 117 |
| 43 | 7159 | 7237 | 847 | 749 | 195 | 163 | 111 |
| 45 | 6781 | 6853 | 808 | 736 | 195 | 163 | 104 |
| 47 | 6026 | 6123 | 775 | 743 | 195 | 150 | 111 |
| 49 | 6000 | 6104 | 801 | 749 | 189 | 150 | 111 |
| 51 | 6000 | 6058 | 795 | 749 | 182 | 143 | 104 |
| 53 | 6006 | 6052 | 808 | 749 | 176 | 143 | 104 |
| 55 | 5987 | 5993 | 808 | 749 | 176 | 143 | 104 |
| 57 | 5921 | 5960 | 801 | 736 | 163 | 143 | 104 |
| 59 | 5876 | 5889 | 801 | 736 | 156 | 143 | 104 |
| 61 | 5804 | 5843 | 795 | 730 | 143 | 130 | 91 |
| 63 | 5765 | 5785 | 762 | 717 | 130 | 124 | 85 |

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Table 3. ${ }^{1} \mathrm{H}-{ }_{-}^{13} \mathrm{C}$ dipolar coupling constants for HBO 1 MHOB

| at different temperatures. |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $T /{ }^{\circ} \mathrm{C}$ | $\alpha-\alpha$ | $\alpha-\beta$ | $\alpha-12$ | $\alpha-11$ | $\alpha-9$ |
| 44 | 4818 | 816 | 127 | 496 | 67 |
| 46 | 4785 | 765 | 130 | 491 | 66 |
| 48 | 4712 | 747 | 126 | 485 | 64 |
| 50 | 4639 | 683 | 115 | 479 | 65 |
| 52 | 4562 | 651 | 107 | 478 | 58 |
| 54 | 4496 | 619 | 101 | 480 | 56 |
| 56 | 4417 | 638 | 95 | 475 | 56 |
| 58 | 4337 | 640 | 90 | 464 | 50 |
| 60 | 4249 | 637 | - | 446 | - |
| 62 | 4176 | 614 | - | 430 | - |
| 64 | 4091 | 634 | - | 449 | - |
| 66 | 3990 | 607 | - | 421 | - |
| 68 | 3900 | 571 | - | 423 | - |
| 70 | 3795 | 540 | - | 407 | - |
| 72 | 3682 | 505 | - | 391 | - |
| 74 | 3553 | 528 | - | 365 | - |
| 76 | 3392 | 474 | - | 353 | - |
| 78 | 3223 | 459 | - | 323 | - |
| 80 | 3010 | 419 | - | 313 | - |
| 82 | 2737 | 348 | - | - | - |
| 84 | 2332 | 319 | - | - | - |

## 3.2. ${ }^{2} \mathrm{H}$ quadrupole coupling

For the monodeuterated compounds, the deuterium quadrupole coupling constants can be determined from the ${ }^{2} \mathrm{H}$ NMR spectra directly. Each of the deuterium NMR spectra of $3 \mathrm{M} 2 \mathrm{C} 2 \mathrm{DPOOB}, 4 \mathrm{M} 2 \mathrm{C} 2 \mathrm{DPOOB}$, and HBO1M1DHOB shows a doublet. The doublet separation is due to deuterium quadrupole splitting, which is related to the components of the order parameter
tensor of the $\mathrm{C}-\mathrm{D}$ bond by

$$
\begin{equation*}
\Delta v=\left(3 e^{2} q Q / 2 h\right)\left[S_{\mathrm{aa}}+\eta\left(S_{\mathrm{bb}}-S_{\mathrm{cc}}\right)\right] . \tag{3}
\end{equation*}
$$

If it is assumed that the asymmetry parameter $\eta$ is zero in equation (3), $S_{\text {aa }}$ is simply the order parameter of the C-D bond, $S_{\mathrm{CD}}$. By using the commonly accepted value of $e^{2} q Q / h=168 \mathrm{kHz}$ for deuterium bonded to aliphatic carbons, the value of $S_{\mathrm{CM}}$ can be readily calculated from $\Delta v$, which is taken to be negative from geometric considerations.

The values of the order parameters $S_{\mathrm{CD}}$ for the three compounds studied are listed in tables 4,5 , and 6 . For comparison, the corresponding $S_{\text {CH }}$ values were calculated from ${ }^{2} \mathrm{H}_{-}{ }^{13} \mathrm{C}$ dipolar coupling constants (tables 1 , 2 , and 3 ) by using equation (2) and the relation [6]

$$
\begin{equation*}
S_{\mathrm{CH}}=-4.407 \times 10^{-5} D_{\mathrm{CH}} \tag{4}
\end{equation*}
$$

Table 4. Order parameters $S_{\mathrm{CD}}$ calculated from quadrupolar splittings and dipolar coupling constants for the $\alpha^{\prime}$ carbon in 4M2C2DPOOB at different temperatures.

| $T /{ }^{\circ} \mathrm{C}$ | $S_{\mathrm{CD}}$ from <br> quadrupolar splittings | $S_{\mathrm{CD}}$ from dipolar <br> coupling constant |
| :---: | :---: | :---: |
| 55 | -0.319 | -0.306 |
| 56 | -0.319 | -0.306 |
| 57 | -0.318 | -0.305 |
| 58 | -0.317 | -0.304 |
| 59 | -0.316 | -0.303 |
| 60 | -0.315 | -0.301 |
| 61 | -0.313 | -0.300 |
| 62 | -0.312 | -0.299 |
| 63 | -0.311 | -0.297 |
| 64 | -0.309 | -0.296 |
| 65 | -0.307 | -0.295 |
| 66 | -0.306 | -0.294 |
| 67 | -0.304 | -0.293 |

Table 5. Order parameters $S_{\text {CD }}$ calculated from quadrupolar splittings and dipolar coupling constants for the $\alpha^{\prime}$ carbon in 3M2C2DPOOB at different temperatures.

| $T /{ }^{\circ} \mathrm{C}$ | $S_{\mathrm{CD}}$ from <br> quadrupolar splittings | $S_{\mathrm{CD}}$ from dipolar <br> coupling constant |  |
| :---: | :---: | :---: | :---: |
| 47 | -0.266 | -0.270 | -0.296 |
| 49 | -0.264 | -0.269 | -0.296 |
| 51 | -0.264 | -0.267 | -0.294 |
| 53 | -0.265 | -0.267 | -0.293 |
| 55 | -0.264 | -0.264 | -0.291 |
| 57 | -0.261 | -0.263 | -0.289 |
| 59 | -0.259 | -0.260 | -0.286 |
| 61 | -0.256 | -0.258 | -0.284 |
| 63 | -0.254 | -0.255 | -0.280 |

Table 6. Order parameters $S_{\mathrm{CD}}$ calculated from quadrupolar splittings and dipolar coupling constants for the $\alpha^{\prime}$ carbon in HBO1M1DHOB at different temperatures.

| $T /{ }^{\circ} \mathrm{C}$ | $S_{\mathrm{CD}}$ from <br> quadrupolar splittings | $S_{\mathrm{CD}}$ from dipolar <br> coupling constant |
| :---: | :---: | :---: |
| 44 | -0.209 | -0.212 |
| 46 | -0.206 | -0.211 |
| 48 | -0.203 | -0.208 |
| 50 | -0.200 | -0.204 |
| 52 | -0.197 | -0.201 |
| 54 | -0.194 | -0.198 |
| 56 | -0.191 | -0.195 |
| 58 | -0.187 | -0.191 |
| 60 | -0.184 | -0.187 |
| 62 | -0.180 | -0.184 |
| 64 | -0.176 | -0.180 |
| 66 | -0.172 | -0.176 |
| 68 | -0.167 | -0.172 |
| 70 | -0.163 | -0.167 |
| 72 | -0.157 | -0.162 |
| 74 | -0.151 | -0.157 |
| 76 | -0.144 | -0.149 |
| 78 | -0.136 | -0.142 |
| 80 | -0.125 | -0.133 |
| 82 | -0.113 | -0.121 |
| 84 | -0.088 | -0.103 |

and the data are listed in the same tables. The small systematic differences are probably due to the assumptions involved in the calculations ( $e^{2} q Q=168 \mathrm{kHz}, \eta=0$, and $\left.r_{\mathrm{CD}}=0.110 \mathrm{~nm}\right)$.

For 3 M 2 C 2 DPOOB and 4 M 2 C 2 DPOOB , the order parameters of the C-D bonds, $S_{\mathrm{CD}}$, become slightly more negative with the decrease of temperature, indicating an increase in the ordering of the smectic phases. Compared with HBO1M1DHOB, the values of $S_{\text {CD }}$ for 3 M 2 C 2 DPOOB and 4 M 2 C 2 DPOOB are substantially larger in absolute value. It is possible that the large chlorine atom bonded to the $\alpha^{\prime}$-carbon causes a considerable change in the $\mathrm{C}-\mathrm{C}-\mathrm{H}$ bond angle, so that the average conformation of the aliphatic chain is quite different from a normal alkyl chain. During the transition from the SmC to SmA phase, there is a discontinuity in the order parameters.

### 3.3. Comparison with $2 D{ }^{13} C$ NMR results

The orientational ordering of the protonated compound 2 M 2 CPOOB was previously studied using the 2D SLF/OMAS method, and all one-bond and some two-bond dipolar coupling constants were determined [9]. Thus, the C-H dipolar coupling constants for the $\alpha^{\prime}$-carbon determined from the 2D method can be compared with that obtained from the 1D spectra of the monodeuterated compound. The data obtained from the two types of experiment at two temperatures are shown in table 7 , and they agree with each other very

Table 7. Comparison of one-bond ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ dipolar couplings $(\mathrm{Hz})$ for the $\alpha^{\prime}$-carbon in 3M2CPOOB obtained from two kinds of experiments. Two sets of diastereo isomers, a $[(2 S, 3 S)$ and $(2 R, 3 R)]$ pair and a $[(2 S, 3 R)$ and $(2 R, 3 S)]$ pair, were present in the compound.

| $T /{ }^{\circ} \mathrm{C}$ | From ${ }^{2} \mathrm{H}^{1{ }^{1}} \mathrm{C}$ | From SLF/OMAS |
| :---: | :---: | :---: |
| 45 (smectic C phase) | $6781 ; 6853$ | 6780 |
| 58 (smectic A phase) | $5921 ; 5960$ | 5880 |

well. Actually, the ${ }^{2} \mathrm{H}^{13} \mathrm{C}$ dipolar coupling constants obtained from the 1D spectra are more accurate, because the ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ dipolar coupling constants obtained from the SLF/OMAS method depend on the accuracy of the scaling factor of the dipolar decoupling sequence and the angle between the spinning axis and the magnetic field.

When optically pure L-isoleucine $[(2 S, 3 S)-(+)-2$-amino3 -methylpentanoic acid] is used as the starting material to synthesize 3M2CPOOB (scheme 1), the product contains only one optical isomer. When racemic dLisoleucine is used as the starting material, because the molecule has two chiral centres, the final product is composed of two sets of diastereo isomers, a [( $2 S, 3 S$ ) and $(2 R, 3 R)]$ pair and a $[(2 S, 3 R)$ and $(2 R, 3 S)]$ pair. In the synthesis of the monodeuterated compound (scheme 2 ), the optical isomers produced in the deuteration step were not isolated, and the final product 3M2C2DPOOB has two sets of diastereo isomers. It is interesting to observe that the two pairs of diastereo isomers show different ${ }^{13} \mathrm{C}$ peaks for the $\alpha^{\prime}, \beta^{\prime}, \omega^{\prime}$, and $\omega^{\prime \prime}$ carbons in the chiral chain of the protonated compound (figure 2, top); separate peaks for the $\gamma^{\prime}$ carbon are not observed because of overlapping. Furthermore, the ${ }^{2} \mathrm{H}^{-13} \mathrm{C}$ dipolar splittings for the two sets of diastereo isomers are also dissimilar (figure 2, bottom and table 2). The different ${ }^{13} \mathrm{C}$ chemical shifts and ${ }^{2} \mathrm{H}-{ }^{13} \mathrm{C}$ dipolar splittings imply that the orientational ordering (conformational averaging) of the $\mathrm{C}-\mathrm{H}$ bonds in the chiral chain of these two sets of diastereo isomers is quite different, but we cannot assign the data to the individual pairs. In contrast, these kinds of difference were too small to be clearly distinguished in the SLF/OMAS experiments [9] for 3M2CPOOB, but were observed for the racemates of $4^{\prime}$-(3-methyl-2-bromopentanoyloxy)-4-hexyloxybiphenyl s [13].
${ }^{13} \mathrm{C}$ chemical shift data, when combined with results obtained from 2D SLF/OMAS experiments, can be used to study the temperature dependence of the order parameters over a wide temperature range [6]. The $S_{\mathrm{CH}}$ values thus calculated for the $\alpha^{\prime}$-carbon in 3 M 2 CPOOB at 15 temperatures are compared with those obtained in this work, and the results are presented in table 8 . The data are very similar except that we can determinate the order parameters of the two diastereoisomers with the

Table 8. Comparison of local order parameters $S_{\mathrm{CH}}$ of the $\alpha^{\prime}$-carbon in 3 M 2 CPOOB obtained from two kinds of experiments. Two sets of diastereo isomers, a $[(2 S, 3 S)$ and $(2 R, 3 R)]$ pair and a $[(2 S, 3 R)$ and $(2 R, 3 S)]$ pair, were present in the compound.

From SLF/OMAS combined

| $T /{ }^{\circ} \mathrm{C}$ | From ${ }^{2} \mathrm{H}^{13} \mathrm{C}$ | with ${ }^{13} \mathrm{C}$ chemical shift |  |
| :---: | :---: | :---: | :---: |
| 35 | $-0.336 ;$ | -0.340 | -0.315 |
| 37 | $-0.331 ;$ | -0.337 | -0.312 |
| 39 | $-0.328 ;$ | -0.332 | -0.300 |
| 41 | $-0.321 ;$ | -0.327 | -0.308 |
| 43 | $-0.315 ;$ | -0.319 | -0.292 |
| 45 | $-0.299 ;$ | -0.302 | -0.285 |
| 47 | $-0.266 ;$ | -0.270 | -0.269 |
| 49 | $-0.264 ;$ | -0.269 | -0.265 |
| 51 | $-0.264 ;$ | -0.267 | -0.265 |
| 53 | $-0.265 ;$ | -0.267 | -0.265 |
| 55 | $-0.264 ;$ | -0.264 | -0.262 |
| 57 | $-0.261 ;$ | -0.263 | -0.260 |
| 59 | $-0.259 ;$ | -0.260 | -0.262 |
| 61 | $-0.256 ;$ | -0.258 | -0.262 |
| 63 | $-0.254 ;$ | -0.255 | -0.256 |

1D experiment. In both sets of data, there is a discontinuity in the order parameters during the transition from the SmC phase to $\operatorname{SmA}$ phase $\left(46^{\circ} \mathrm{C}\right)$. This discontinuity was observed previously [11] and was attributed to the first order nature of the $\mathrm{SmC}-\mathrm{SmA}$ transition for this liquid crystal.

### 3.4. Temperature dependence

The temperature dependence of the order parameters over a wide temperature range was studied for the liquid crystal 3 M 2 CPOOB with the combination of PELF/OMAS results and ${ }^{13} \mathrm{C}$ chemical shift data, which were obtained from one-dimensional NMR experiments. The data for the aromatic rings and for the aliphatics carbons are presented in figures 4 and 5, respectively. They are fitted to the Haller equation:

$$
\begin{equation*}
S(T)=S_{\mathrm{o}}\left(1-T / T^{*}\right)^{F} \tag{5}
\end{equation*}
$$

where $S_{\mathrm{o}}$ and $F$ are empirical constants and $T^{*}$ is a temperature at which the order parameter becomes zero.

The values of $S_{\mathrm{o}}$ and $F$ obtained by fitting the order parameters calculated from chemical shifts to equation (5) are listed in table $9 . T^{*}$ was initially treated as a variable parameter for fitting each curve; then, the average value


Figure 4. Plot of order parameters of the aromatic rings of 3 M 2 CPOOB against $T / T^{*}$, with $T^{*}=336.8 \mathrm{~K}$. The solid line was calculated from equation (5).


Figure 5. Plot of $\mathrm{C}-\mathrm{H}$ order parameters of the aliphatic carbons of 3M2CPOOB against $T / T^{*}$, with $T^{*}=336.8 \mathrm{~K}$. The symbols are: $\bullet=\alpha^{\prime}, \bullet=\beta^{\prime}, \diamond=\alpha,+=\xi, \times=\eta$, and $\Delta=\omega$. The solid lines were calculated from equation (5).
of 336.8 K was used as a fixed parameter for all final fittings. As expected, the values of $S_{\mathrm{o}}$ for the phenyl rings in 3M2CPOOB are close to unity (table 9). For a perfectly aligned all-trans-conformation of the aliphatic chain, $S_{\mathrm{o}}$ for each C-H bond should be close to -0.333 [i.e. $\left.\left(3 \cos ^{2} 109^{\circ} 28^{\prime}-1\right) / 2\right]$. Only two of the values listed

Table 9. Values of $S_{\mathrm{o}}$ and $F$ for 3M2CPOOB obtained by fitting the calculated order parameters to equation (5); $R$ is the correlation coefficient.

| C | Ring | $\alpha^{\prime}$ | $\beta^{\prime}$ | $\alpha$ | $\xi$ | $\eta$ | $\omega$ |
| :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: |
| $S_{\mathrm{o}}$ | 1.04 | -0.411 | -0.300 | -0.376 | -0.222 | -0.116 | -0.099 |
| $F$ | 0.15 | 0.14 | 0.15 | 0.14 | 0.09 | 0.07 | 0.09 |
| $R$ | 0.995 | 0.991 | 0.965 | 0.991 | 0.975 | 0.988 | 0.993 |

in table 9 (the $\alpha$ and $\beta^{\prime}$ atoms) are close to this estimation. The value for the $\alpha^{\prime}$ carbon is more negative because it is connected to a carboxy segment, and the increased mobility towards the end of the aliphatic chain reduces the limiting order parameter $S_{0}$ for the three outer carbons atom ( $\xi, \eta$ and $\omega$ ).

Although the long range dipolar coupling constants cannot be used directly to calculate the order parameters, they are also determined by the orientational ordering of the liquid crystal molecules and may obey a relation similar to the Haller equation, equation (6). Therefore, we tried to fit the data listed in table 2 to the following equation:

$$
\begin{equation*}
D(T)=D_{\mathrm{o}}\left(1-T / T^{*}\right)^{F} \tag{6}
\end{equation*}
$$

where $D_{\mathrm{o}}$ and $F$ are empirical constants, and $T^{*}$ is a temperature at which the dipolar coupling constants become zero and should be the same as that in equation (5). The results are presented in figure 6, and the values of $D_{\text {o }}$ and $F$ are listed in table 10 . The values of $F$ are relatively constant, but the physical significance is not clear.


Figure 6. Plot of long range dipolar coupling constants of the carbons of 3M2CPOOB against $T / T^{*}$, with $T^{*}=336.8 \mathrm{~K}$. The symbols are: $\bullet$ and $=\alpha^{\prime}-\beta^{\prime}, \diamond$ and $\bigcirc=\alpha^{\prime}-\omega^{\prime \prime}$, and $+=\alpha^{\prime}-\omega^{\prime}$. The solid lines were calculated from equation (6).

Table 10. Values of $D_{0}$ and $F$ for 3M2CPOOB obtained by fitting the ${ }^{1} \mathrm{H}-{ }^{13} \mathrm{C}$ dipolar coupling constants to equation (6): $R$ is the correlation coefficient.

| C | $\alpha^{\prime}-\beta^{\prime}$ | $\alpha^{\prime}-\beta^{\prime}$ | $\alpha^{\prime}-\omega^{\prime \prime}$ | $\alpha^{\prime}-\omega^{\prime \prime}$ | $\alpha^{\prime}-\omega^{\prime}$ |
| :--- | ---: | ---: | ---: | ---: | ---: |
| $D_{\circ}$ | 956 | 856 | 275 | 221 | 137 |
| $F$ | 0.143 | 0.135 | 0.128 | 0.110 | 0.077 |
| $R$ | 0.987 | 0.992 | 0.994 | 0.977 | 0.990 |

## 4. Conclusion

Long range dipolar coupling constants are important in the investigation of detailed chain motions and rapid rotations or jumps of the rings. The selective deuteration of liquid crystals provides a simple and powerful method for the determination of long range dipolar coupling constants. Ferroelectric liquid crystals have been widely studied because of their interest in display applications. In this report, we have presented an example of using onedimensional ${ }^{13} \mathrm{C}$ NMR to obtain long range $\mathrm{H}-\mathrm{C}$ dipolar couplings in three FLCs. These data can be used to study molecular motions and conformational properties of FLCs over the whole mesomorphic temperature range. The results have been compared with those obtained from the 2D SLF/OMAS method. Additional data for the monodeuterated liquid crystals were obtained from $1 D^{2} \mathrm{H}$ NMR.

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## References

[1] Meyer, R. B., Liebert, L., Strzelecki, L., and Keller, P., 1975, J. Physique (Lett.), 36, L69.
[2] Sakurai, T., Mikami, N., Higuchi, R., Honma, M., Ozaki, M., and Yoshino, K., 1986, J. chem. Soc., chem. Comm., 978.
[3] Bahr, C., and Heppke, G., 1986, Mol. Cryst. liq. Cryst. Lett., 4, 31.
[4] Sierra, T., Serrano, J. L., Ros, M. B., Ezcurra, A., and Zubia, J., 1992, J. Am. chem. Soc., 114, 7645.
[5] Emsley, J. W. (Editor), 1985, Nuclear Magnetic Resonance of Liquid Crystals (Dordrecht: Kluwer).
[6] Fung, B. M., 1996, in Encyclopedia of Nuclear Magnetic Resonance, edited by D. M. Grant and R. K. Harris (Chichester: Wiley), pp. 2744-2751.
[7] Schmidt-Rohr, K., Nanz, D., Emsley, L., and Pines, A., 1994, J. phys. Chem., 98, 6668; Hong, M., SchmidtRohr, K., and Pines, A., 1995, J. Am. chem. Soc., 117, 3310; Hong, M., Pines, A., and Caldarelli, S., 1996, J. phys. Chem., 100, 14 815; Caldarelli, S., Hong, M., Emsley, L., and Pines, A., 1996, J. phys. Chem., 100, 18696.
[8] Fung, B. M., Ermolaev, K., and Yu, Y., 1999, J. magn. Reson., 138, 28.
[9] Poon, C.-D., and Fung, B. M., 1989, Liq. Cryst., 5, 1159; Poon, C.-D., and Fung, B. M., 1989, J. chem. Phys., 91, 7392.
[10] Richter, W., Reimer, D., Fung, B. M., Twieg, R. J., and Betterton, K., 1990, Liq. Cryst., 8, 687.
[11] Ho, M.-H., Fung, B. M., Wand, M., and Vohra, R. T., 1993, Ferroelectrics, 138, 51.
[12] Loubser, C., Wessels, P. L., Goodby, J. W., and Styring, P., 1993, Liq. Cryst., 15, 233.
[13] Chen, A., Poon, C.-D., Dingemans, T. J., and Samulski, E. T., 1998, Liq. Cryst., 24, 255.
[14] Höhener, A., Müller, L., and Ernst, R. R., 1979, Mol. Phys., 38, 909.
[15] Sinton, S. W., Zax, D. B., Murdoch, J. B., and Pines, A., 1984, Mol. Phys., 53, 333.
[16] Sandstrom, D., and Levitt, M. H., 1996, J. Am. chem. Soc., 118, 6966.
[17] Ciampi, E., de Luca, G., and Emsley, J. W., 1997, J. magn. Reson., 129, 207.
[18] Caldarelli, S., Lesage, A., and Emsley, L., 1996, J. Am. chem. Soc., 118, 12 224; Auger, C., Lesage, A., Caldarelli, S., Hodgkinson, P., and Emsley, L., 1998, J. phys. Chem. B, 102, 3718.
[19] Canlet, C., and Fung, B. M., 2000, J. chem. Phys. B, 104, 6181.
[20] Fung, B. M., Ermolaev, K., and Khitrin, A. K., 2000, J. magn. Reson., 142, 97.


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